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## Abstract

**The entire work of this thesis has been divided into five chapters**

Throughout the studies we have introduced a constant  $U$  in the Lagrangian ( $L$ ) in such a way that the energy constant ( $h$ ) vanishes at  $L_4$  (the liberation point). We have used **mobile coordinate** system ( $Q = (x, y)^T$  on the orbit where the modulus of momentary velocity  $V(t) = |\dot{Q}(t)| = \sqrt{\dot{x}^2 + \dot{y}^2} \neq 0$ ) to determine the periodic orbits by giving displacement to these coordinates along the **normal** ( $N$ ) and the **tangent** ( $M$ ) directions. We have constructed an algorithm, in two stages, to draw the periodic orbits. These are: first predictor-part and then corrector-part. In each chapter, we have drawn six (or five) families of periodic orbits. And in each family, we have drawn five figures corresponding to the different values of  $h$ . These orbits have been numbered 1, 2, 3, 4 and 5 corresponding to values of  $h$  mentioned in each figure on the left hand top of each figure in each chapter. It is observed that the final orbit passing through the libration point  $L_4$ , in each case, is non-symmetrical and therefore, the family can be further continued whereas in the case of Karimov and Sokolsky (1989) model, family terminates when the orbit touches the point  $L_4$ .

The entire work of this thesis has been divided in five chapters.

The chapter-1 is introductory in nature. It contains history and development of the problem.

In all other chapters (2 to 5), we have drawn the periodic orbits around the triangular libration point  $L_4$ , in the restricted three body problem when the primaries are axis symmetric rigid bodies with radiation pressure. The equatorial plane of the oblate body of mass  $m_2$  is coincident with the plane of motion. All the chapters are divided into 8 sections.

In each chapters, we have drawn periodic orbits for the following:

For chapter-2 (i) for fixed  $\mu = .001$ ,  $A = 0.0$  (Fig 1), (ii) for fixed  $\mu = .001$ ,  $A = .01$  (Fig 2), (iii) for fixed  $\mu = .001$ ,  $A = .001$  (Fig 3), (iv) for fixed  $\mu = .001$ ,  $A = .0001$  (Fig 4), (v) for fixed  $\mu = .001$ ,  $A = .00001$  (Fig 5), (vi) for fixed  $\mu = .01$ ,  $A = .001$  (Fig 6).

For chapter-3 (i) for fixed  $\mu = .001$ ,  $\sigma_1 = 0.0$  and  $\sigma_2 = .001$  (Fig 1), (ii) for fixed  $\mu = .001$ ,  $\sigma_1 = .0001$  and  $\sigma_2 = .001$  (Fig 2), (iii) for fixed  $\mu = .001$ ,  $\sigma_1 = .001$  and  $\sigma_2 = .001$  (Fig 3), (iv) for fixed  $\mu = .001$ ,  $\sigma_1 = .001$  and  $\sigma_2 = .002$  (Fig 4), (v) for fixed  $\mu = .001$ ,  $\sigma_1 = .002$  and  $\sigma_2 = .003$  (Fig 5).

For chapter-4 (i) for fixed  $\mu = .001$ ,  $A_1 = 0.0$ ,  $A_2 = 0.0$ ,  $A'_1 = 0.001$  and  $A'_2 = 0.0$  (Fig 1), (ii) for fixed  $\mu = .001$ ,  $A_1 = .001$ ,  $A_2 = 0.0$ ,  $A'_1 = .001$  and  $A'_2 = 0.0$  (Fig 2), (iii) for fixed  $\mu = .001$ ,  $A_1 = .001$ ,  $A_2 = .001$ ,  $A'_1 = 0.0$  and  $A'_2 = 0.0$  (Fig 3), (iv) for fixed  $\mu = .001$ ,  $A_1 = .001$ ,  $A_2 = .001$ ,  $A'_1 = .001$  and  $A'_2 = .001$  (Fig 4), (v) for fixed  $\mu = .001$ ,  $A_1 = 0.002$ ,  $A_2 = .003$ ,  $A'_1 = .004$  and  $A'_2 = .005$  (Fig 5).

For chapter-5 (i) for fixed  $\mu = .001$ ,  $A = .001$  and  $p = 0.0$  (Fig 1), (ii) for fixed  $\mu = .001$ ,  $A = .001$  and  $p = .0001$  (Fig 2), (iii) for fixed  $\mu = .001$ ,  $A = .001$  and  $p = 0.001$  (Fig 3), (iv) for fixed  $\mu = .001$ ,  $A = .001$  and  $p = 0.01$  (Fig 4), (v) for fixed  $\mu = .001$ ,  $A = .001$  and  $p = 0.1$  (Fig 5).

It has been observed that by taking axis symmetric rigid bodies with radiation pressure, the families of periodic orbits continues beyond  $L_4$  whereas in case of Karimov and Sokolsky (1989), who have not taken axis symmetric rigid bodies, the families terminates at  $L_4$ .